

Quantum Mechanics in one space General Relativity in another

Connor Behan

November 28, 2025
Acadia University

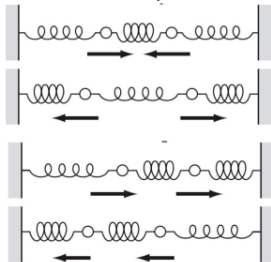


A picture to have in mind

Goal: Map a seemingly complicated system onto a simple one.

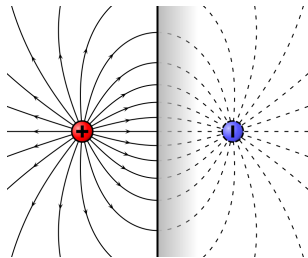
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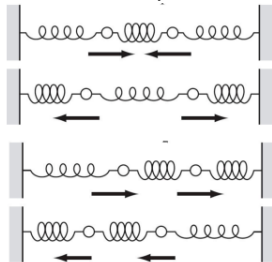
← Normal modes

Image charges →



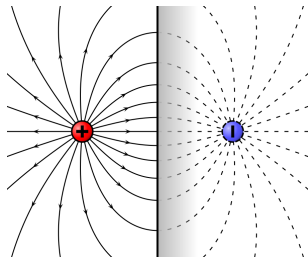
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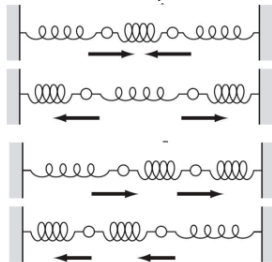
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An intriguing example is the **holographic principle**.

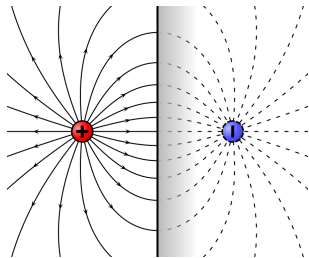
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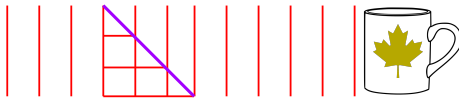


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Image charges →



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← 3d scene (complicated)



← 2d plate (simple)

Gravity vs other forces

Maxwell's eq'ns determine **fields on** spacetime.

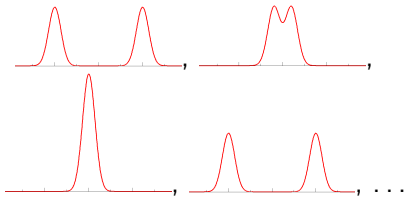
$$\begin{aligned}\nabla \cdot \mathbf{E} &= 4\pi\rho, & \nabla \times \mathbf{E} &= -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} \\ \nabla \cdot \mathbf{B} &= 0, & \nabla \times \mathbf{B} &= \frac{1}{c} \left(4\pi\mathbf{J} + \frac{\partial \mathbf{E}}{\partial t} \right)\end{aligned}$$

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Maxwell's eq'ns determine **fields on** spacetime. They are linear.

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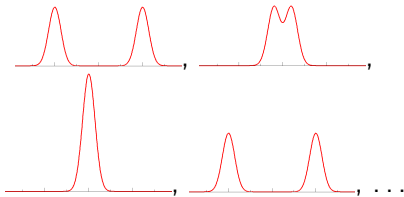


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Einstein's equation determines the **shape of** spacetime

$$g_{\mu\nu} = \text{diag}(-1, 1, 1, 1) \Rightarrow ds^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2$$

$$\text{Any } g_{\mu\nu} \Rightarrow ds^2 = \sum_{\mu, \nu=0}^3 g_{\mu\nu} dx^\mu dx^\nu$$

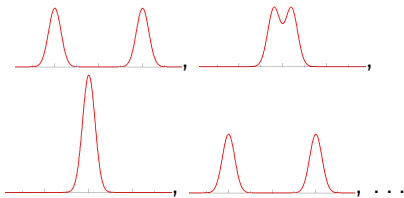
and particles follow geodesics which minimize ds .

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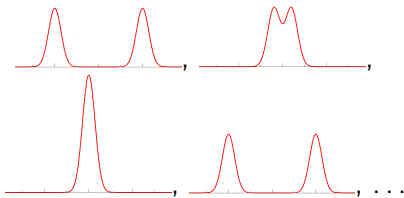
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A simple but profound solution

“Point mass” solution found in [\[Schwarzschild, 1916\]](#) :

$$ds^2 = - \left(1 - \frac{2GM}{c^2 r} \right) c^2 dt^2 + \left(1 - \frac{2GM}{c^2 r} \right)^{-1} dr^2 + r^2 d\Omega^2.$$

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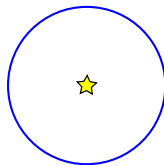
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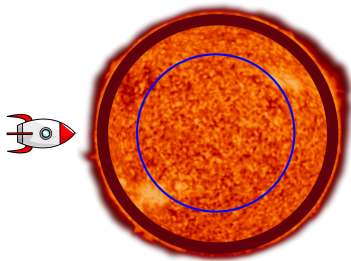


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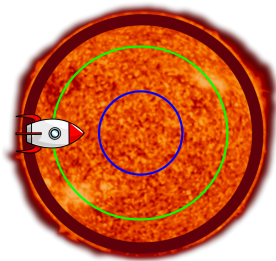


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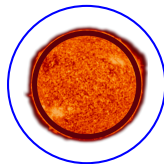


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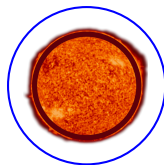
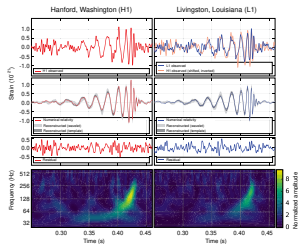


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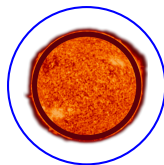
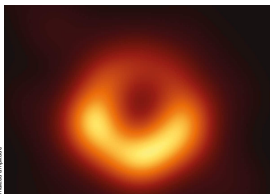
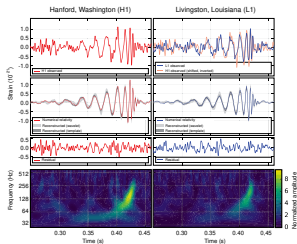
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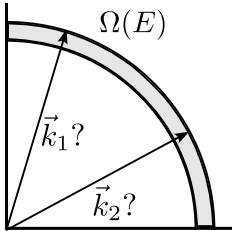
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Most efficient packing of mass!

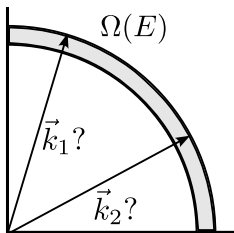
How much information lies inside?

Event horizons also have an interesting entropy [\[Bekenstein, 1973\]](#) .



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$$S = k \log \Omega$$

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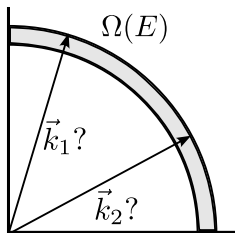
Access to both photons: $S = 0$

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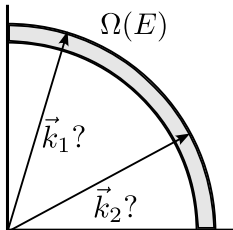
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Entangling correlations should happen across the area: $S \propto A$.

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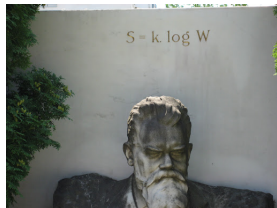


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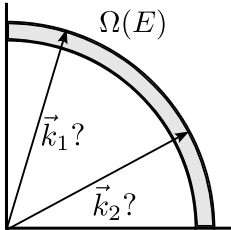
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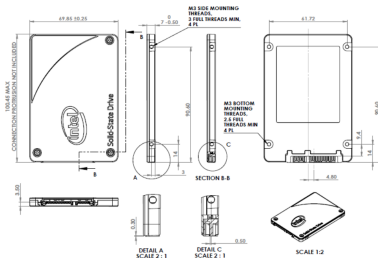
$$A = 0.0163\text{m}^2$$

Current information using E&M: 1TB

Maximum information using E&M: ∞

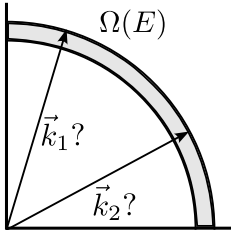
Maximum information using gravity:

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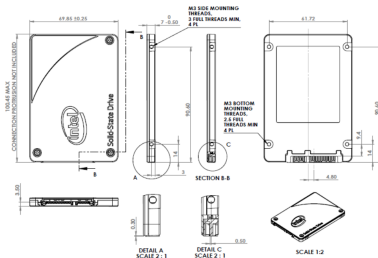
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Most efficient packing again!

Hawking's derivation of the entropy

Thermal state combines $S = k \log \Omega$ and $dE = TdS$:

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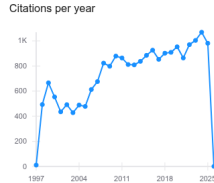
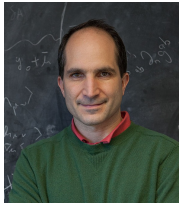
After a black hole has formed, there is a **different vacuum** killed by “negative frequency” operators.

$$b_i = \sum_j \alpha_{ij} a_j + \beta_{ij} a_j^\dagger, \quad N_i = {}_a\langle 0 | b_i^\dagger b_i | 0 \rangle_a = \sum_j |\beta_{ij}|^2.$$

Coefficients found by projecting ingoing wave with constant $v = ct + r$ onto outgoing wave with constant $u = ct - r$ [Hawking, 1975] .

$$u = -\frac{4GM}{c^2} \log(v - v_0) + \dots \Rightarrow kT = \frac{\hbar c^3}{8\pi GM}$$

A well defined holographic duality



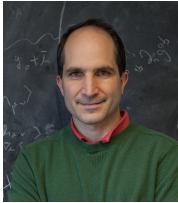
[Maldacena; [hep-th/9711200](#)]

Quantum Field Theory

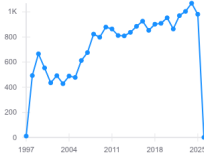


Quantum Gravity

A well defined holographic duality



Citations per year



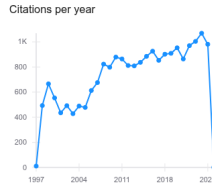
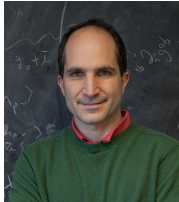
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Conformally invariant
Quantum Field Theory



Asymptotically Anti-de Sitter
Quantum Gravity

A well defined holographic duality



[Maldacena; hep-th/9711200]

Conformally invariant
Quantum Field Theory



Asymptotically Anti-de Sitter
Quantum Gravity

This **AdS/CFT correspondence** (discovered using string theory) has implications for the following questions.

1. How to make the extra dimensions of string theory compact?
2. What does string theory predict non-perturbatively?
3. How to treat strings (even perturbatively) when the non-compact space is curved?
4. What changes when a QFT becomes strongly coupled?

What is AdS?

$$R_{\mu\nu} - \frac{1}{2}R_{\rho\sigma}g^{\rho\sigma}g_{\mu\nu} + \Lambda g_{\mu\nu} = 0$$

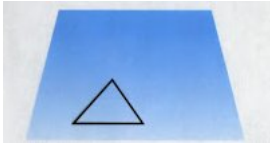
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$\Lambda > 0$, de Sitter



$\Lambda = 0$, Minkowski



$\Lambda < 0$, Anti-de Sitter



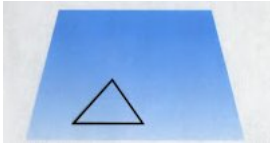
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$$ds^2 = - \left(1 + \frac{r^2}{L^2}\right) dt^2 + \left(1 + \frac{r^2}{L^2}\right)^{-1} dr^2 + r^2 d\Omega^2, \quad \Lambda = \frac{d(1-d)}{2L^2} \quad \begin{array}{c} \uparrow \\ \text{---} \end{array}$$

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 $SO(d+1, 1)$



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 $SO(d, 1) \times \text{Translations}$



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$P = -\rho$ equivalent to
 $\Lambda > 0$ is observed??

[DESI, 2024]

What is AdS?

$$R_{\mu\nu} - \frac{1}{2}R_{\rho\sigma}g^{\rho\sigma}g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi GT_{\mu\nu} = \text{diag}(\rho, P, P, P)$$

$\Lambda > 0$, de Sitter
 $SO(d+1, 1)$



$\Lambda = 0$, Minkowski
 $SO(d, 1) \times \text{Translations}$



$\Lambda < 0$, Anti-de Sitter
 $SO(d, 2)$



$$ds^2 = - \left(1 + \frac{r^2}{L^2}\right) dt^2 + \left(1 + \frac{r^2}{L^2}\right)^{-1} dr^2 + r^2 d\Omega^2, \quad \Lambda = \frac{d(1-d)}{2L^2} \quad \nearrow$$



$P = -\rho$ equivalent to
 $\Lambda > 0$ is observed??

[DESI, 2024]

Dark energy: mysterious cosmic force appears to be weakening, say scientists

Findings could open up possibility the universe will end in a reverse Big Bang or Big Crunch. It's not cosmologists' favorite scenario.

Explainer: Is dark energy destined to dominate the universe?



Long-exposure photos taken at the Dark Energy Survey (DES) National Observatory, where astronomers are studying the nature of dark energy. Photograph: AFP

What is CFT?

Translations	$x'_\mu = x_\mu + a_\mu$	d
Rotations	$x'_\mu = \Lambda_\mu{}^\nu x_\nu$	$\binom{d}{2}$
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Special	$x'_\mu = \frac{x_\mu - b_\mu x^2}{1 - 2b \cdot x + b^2 x^2}$	d

$SO(d, 2)$ symmetric QFT
with **no preferred scale**.

Realized in nature in phase
transitions with $\xi = \infty$.

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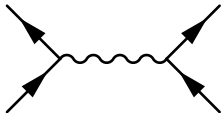
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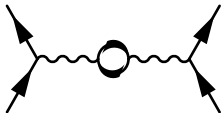


Diagram F is a divergent function of e_0 .

In turn, make e_0 a divergent function of e to
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Resulting $F(e(\mu), \mu)$ must be μ -independent.

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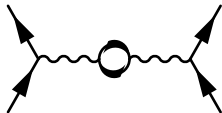
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[Kos, Poland, Simmons-Duffin; 1406.4858]

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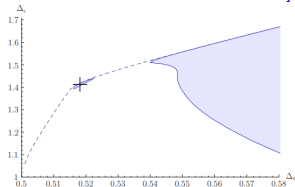


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How do they fit together?

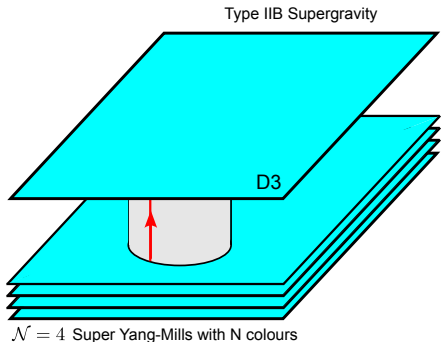
Amplitudes in terms of $s = p_1 \cdot p_2$, $t = p_2 \cdot p_3$, $u = p_2 \cdot p_4$:

$$\textbf{Open} : \frac{\Gamma(-\alpha' s)\Gamma(-\alpha' t)}{\Gamma(1 + \alpha' u)}, \quad \textbf{Closed} : \frac{\Gamma(-\alpha' s)\Gamma(-\alpha' t)\Gamma(-\alpha' u)}{\Gamma(1 + \alpha' s)\Gamma(1 + \alpha' t)\Gamma(1 + \alpha' u)}$$

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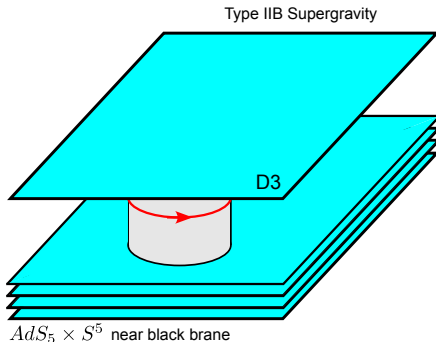
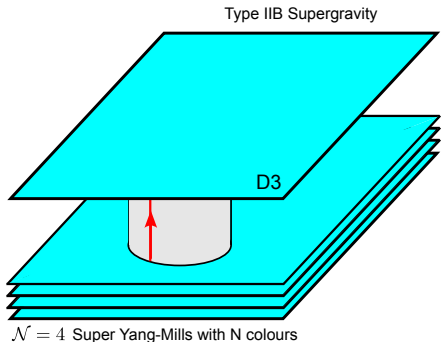
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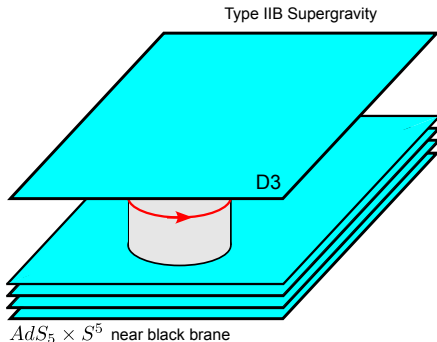
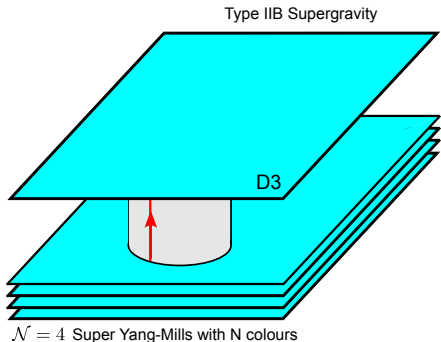
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Type IIB string theory on $AdS_5 \times S^5$ is equivalent to a CFT labelled by N and g_{YM} [Maldacena; hep-th/9711200] .

Identifying parameters

Path integrate over manifolds, not just metrics:

$$A = \sum_{top} \int DX Dg e^{-S - \lambda S_{top}} \dots$$

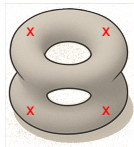
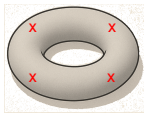
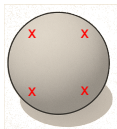
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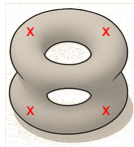
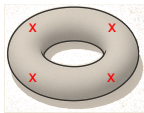
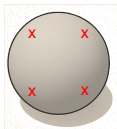


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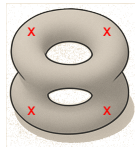
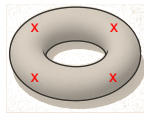
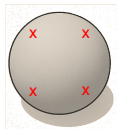
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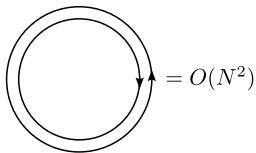
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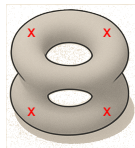
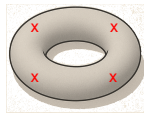
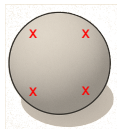


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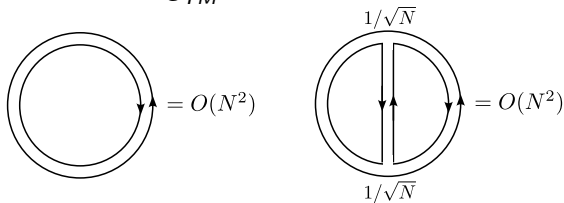
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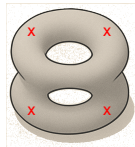
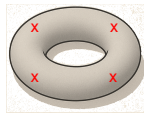
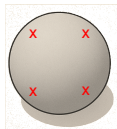


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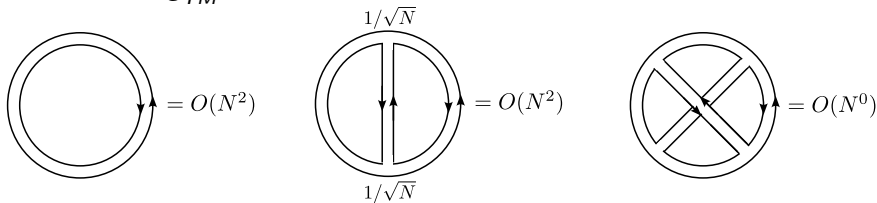
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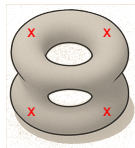
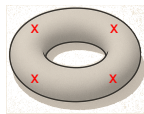
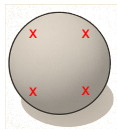


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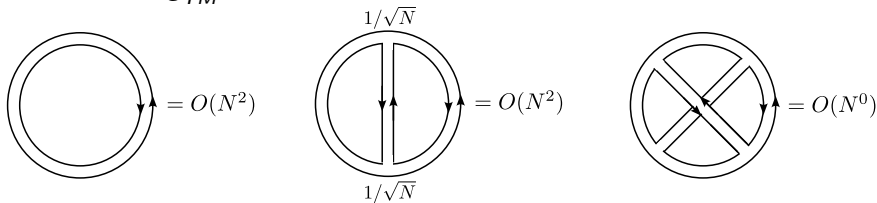
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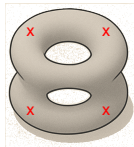
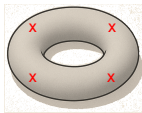
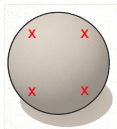
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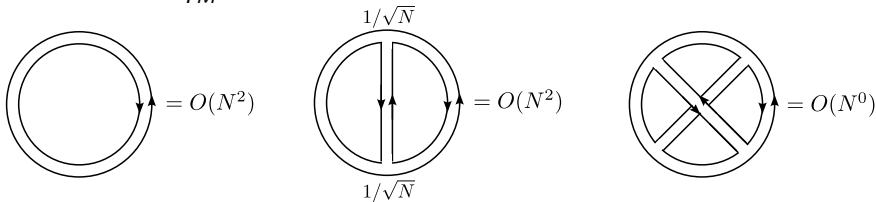
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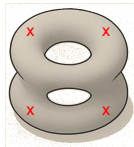
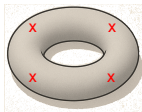
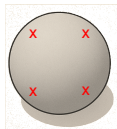
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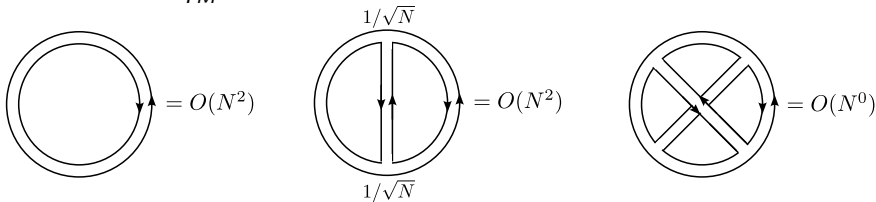
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Calculations in AdS are most tractable with $L/l_p \ll 1$, $L/l_s \ll 1$ (large N , large λ)!

Some basic checks

Asymptotic density of states $S = \log \Omega$ can be computed by dimensional analysis in a CFT.

$$E \propto VT^{d+1}, \quad S \propto VT^d \quad \Rightarrow \quad S \propto E^{\frac{d}{d+1}}$$

$\mathcal{N} = 4$ Super Yang-Mills is $3 + 1$ dimensional so $S = E^{3/4}$.

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High energy states in $AdS_5 \times S^5$ are large black holes.

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Some gauge theory phases at lower energy are harder to study without the correspondence [\[Witten; hep-th/9803131\]](#).

Beyond counting operators

Consider low energy (or really dimension) operators.

$$\mathcal{O}_{\mu_1 \dots \mu_\ell}(x) \mapsto \lambda^{-\Delta} \mathcal{O}_{\mu_1 \dots \mu_\ell}(\lambda x), \quad (mL)^2 = \Delta(\Delta - d) + \ell(\ell + d - 2)$$

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$T_{\mu\nu} = \phi^i \partial_\mu \partial_\nu \phi^i - 2 \partial_\mu \phi^i \partial_\nu \phi^i - \text{trace}$	Graviton field
...	...

Protected operators with $\mathcal{N} = 4$ SUSY all have corresponding fields in $AdS_5 \times S^5$ gravity [Dolan, Osborn; [hep-th/0209056](#)].

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Protected operators with $\mathcal{N} = 4$ SUSY all have corresponding fields in $AdS_5 \times S^5$ gravity [\[Dolan, Osborn; hep-th/0209056\]](#) .

State-of-the-art SUSY techniques can find terms in $\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle$ **purely from QFT** away from $\lambda \rightarrow \infty$ limit [\[Chester, Pufu; 2003.06412\]](#) . Result matches closed strings for $L \rightarrow \infty$.

Beyond counting operators

Consider low energy (or really dimension) operators.

$$\mathcal{O}_{\mu_1 \dots \mu_\ell}(x) \mapsto \lambda^{-\Delta} \mathcal{O}_{\mu_1 \dots \mu_\ell}(\lambda x), \quad (mL)^2 = \Delta(\Delta - d) + \ell(\ell + d - 2)$$

$J_\mu^{ij} = \phi^i \partial_\mu \phi^j - \phi^j \partial_\mu \phi^i$	Gauge field
$T_{\mu\nu} = \phi^i \partial_\mu \partial_\nu \phi^i - 2 \partial_\mu \phi^i \partial_\nu \phi^i - \text{trace}$	Graviton field
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The same techniques work for another holographic setup with $\mathcal{N} = 2$ SUSY [CB, Chester, Ferrero; 2305.01016] . Result matches open strings for $L \rightarrow \infty$.

Towards more realistic systems

Bound states in QCD obey wave equations which have an AdS_5 interpretation [\[Forshaw, Sandapen; 2501.00526\]](#) .

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Can also model $\mathcal{N} = 4$ Super Yang-Mills when deviations from thermal equilibrium are slowly varying

[Kovtun, Son, Starinets; hep-th/0405241] .

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

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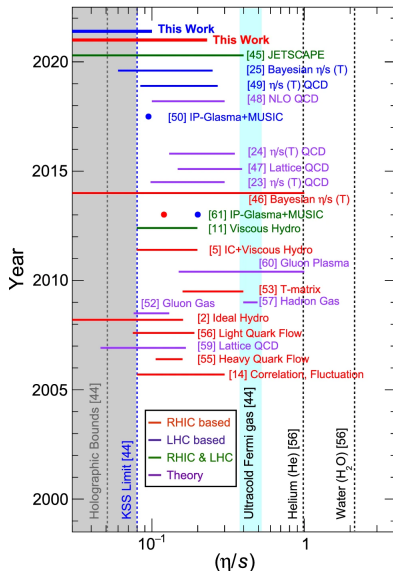
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[Gonzalez, Basu, Guevara, Marin, Pan, Pruneau; 2012.10542]



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Thanks for your attention!